Quantum Uncertainty in Doubly Special Relativity

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The modification of the quantum mechanical commutators in a relativistic theory with an invariant length scale (DSR) is identified. Two examples are discussed where a classical behavior is approached in one case when the energy approaches the inverse of the invariant length which appears as a cutoff in the energy and in the second case when the mass is much larger than the inverse of the invariant length.

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I. INTRODUCTION

One of the basic open problems in theoretical physics is to combine in a consistent way the classical description of the gravitational interaction (general relativity) with quantum mechanics (QM). The analysis of the problems that one finds in different attempts to combine these two theories can be used as a guiding principle to the search of a fundamental theory of quantum gravity.

On the other hand one has direct proposals for this theory like string(M)-theory or loop quantum gravity. But we are not able, presently to establish if any of these or another future proposal is the correct theory. An alternative is to try to identify basic ingredients that this theory should contain and whose details could be a criteria in the future to select among different alternatives. One idea - that has been discussed very often in this context [1] - is the possible appearance of a fundamental length, a scale associated to gravity not just as a dynamical scale but at the kinematical level.

Several arguments going from theoretical analysis in string theory [2] to more or less sophisticated gedanken experiments to measure lengths [1], lead to the conclusion that there is a new contribution to the quantum uncertainty of gravitational origin leading to a length scale as a minimal uncertainty in the determination of space-time coordinates. Some attempts to identify a modification of the quantum mechanical commutators [3] which reproduce the generalized uncertainty principle, have been considered as a way to find one of the ingredients of the quantum theory of gravity.

Another candidate for a signal of quantum gravity effects is the modification of Lorentz symmetry at very high energies, an idea explored intensively in the last years both from a theoretical as well as from a phenomenological perspective [4].

Alternatively, the limitation due to gravity to explore beyond a minimal length, has motivated to consider the possibility of a generalization of the relativity principle compatible with an invariant length scale which is called doubly special relativity (DSR) [5, 6] as a possible ingredient of the flat space-time limit of the quantum theory of gravity. It is remarkable that one can have a generalized relativity principle which is continuously connected with standard Einstein special relativity (SR), compatible with all the tests of SR and leading to corrections accessible to present or near future experiments or even already observed in the high energy tail of the cosmic ray spectrum [7, 8]. Different proposals for DSR theories either based on a deformation of Poincare algebra or by considering directly a modification of the boosts that connect inertial frames has been considered recently from different points of view [5, 6, 9].

It is still not clear how to introduce space-time in DSR [10]. The aim of this paper is to study the modification of QM at the level of commutation relations associated to a modified relativity principle based on a nonlinear realization of Lorentz transformations in energy-momentum space (which is a common ingredient of the different proposals of DSR theories) together with a simple prescription for the space-time sector. We conjecture that the modifications of QM obtained in this way are a remnant of the fundamental theory of QG in the flat space-time limit.

II. MODIFIED QUANTUM MECHANICAL COMMUTATORS

A nonlinear realization of Lorentz transformations in energy-momentum \((E, p)\) space parametrized by an invariant length \(\ell\) can be defined [11] by the relations

\[ \epsilon = E f (\ell E, \ell^2 p^2) \]  \hspace{1cm} (1)

\[ \pi_i = p_i g (\ell E, \ell^2 p^2) \]  \hspace{1cm} (2)

where \((\epsilon, \pi)\) are auxiliary linearly transforming variables which define the nonlinear Lorentz transformation of the
which will be tions appear at the level of the phase space commutators which are still trivial but all the modific-
tivity principle at the level of the space-time coordinate case one does not have any signal of the modi-
cations, one has a linear system of equations for the partial de-
defining the nonlinear realization of Lorentz transform-
ations of special relativity in the limit \( \ell E \ll 1 \). In this

In order to have a quantum theory with such a de-
formed relativity principle, one should find the appro-
piate deformation of relativistic quantum field theory (QFT). First attempts in this direction, based on the
possible connection between a generalization of the rela-
tivity principle and a non-commutativity of space-time, suggesting the formulation of QFT in a non-commutative
space (\( \kappa \)-Minkowsky) as the appropriate deformation of QFT have been explored [12]. But there are general ar-

guments that there will be difficulties to find a realization of a deformed relativity principle along these lines in the

generality idea that in the presence of quantum gravity the quantum mechanical uncertainty principle, and then the
phase space commutators on which it is based, should be modified by terms depending on an invariant length
[3]. Instead of guessing the general structure of the gen-
alized uncertainty principle and the modification of the
phase space commutators leading to such a generalization of
the uncertainty principle, a modification of the con-
mutators (7-10) is obtained directly from a non linear re-
alization of Lorentz transformations in momentum space
parametrized by the functions \( f, g \). In order to discuss the
consequences of the modifications of the quantum me-
chanical commutators one has to specify the non-linear
transformations of energy-momentum. We discuss some
cases in next two sections.

III. AN EXAMPLE WITH AN ENERGY CUTOFF: DSR2

When the two functions \( f, g \) parametrizing the non linear
Lorentz transformations are independent of the mo-
menta (i.e., when \( \partial_2 f = \partial_2 g = 0 \)) the modified quantum

mechanical commutators in (7-10) take a much simpler form

By considering the derivatives with respect to the aux-
iliary energy-momentum variables of the equations (1-2)


defining the nonlinear realization of Lorentz transforma-
tions, one has a linear system of equations for the partial
derivatives required to calculate the phase space com-
mutators (3-6). A straightforward algebra leads to the
modified quantum mechanical commutators

\[
[t, E] = i\hbar \frac{\partial E}{\partial \epsilon} \quad (3)
\]

\[
[t, p_i] = i\hbar \frac{\partial p_i}{\partial \epsilon} \quad (4)
\]

\[
[x_i, E] = i\hbar \frac{\partial E}{\partial \xi_i} \quad (5)
\]

\[
[x_i, p_j] = i\hbar \frac{\partial p_j}{\partial \xi_i}. \quad (6)
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This result can be seen as an explicit realization of a
general idea that in the presence of quantum gravity the
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\[
[t, E] = i\hbar \frac{1}{f + \ell E \partial_1 f} \quad (13)
\]

\[
[t, p_i] = -i\hbar \frac{\partial p_i}{g [f + \ell E \partial_1 f]} \quad (14)
\]

\[
[x_i, E] = 0 \quad (15)
\]

\[
[x_i, p_j] = i\hbar \delta_{ij} \frac{1}{g}. \quad (16)
\]

A very simple choice for the functions \( f, g \)

\[
f = g = (1 - \ell E)^{-1} \quad (17)
\]
is what is known as DSR2 [6] and corresponds to the simplest realization of DSR with an energy cutoff \((E < 1/\ell)\) identified as the inverse of the invariant length. The combinations of derivatives which appear in the phase space commutators are given by

\[
\varpartial_1 g = f + \ell E \varpartial_1 f = (1 - \ell E)^{-2}
\]

and then one has

\[
[t, E] = i\hbar (1 - \ell E)^2
\]

\[
[t, p_i] = -i\hbar \ell p_i (1 - \ell E)
\]

\[
[x_i, E] = 0
\]

\[
[x_i, p_j] = i\hbar \delta_{ij} (1 - \ell E)
\]

a result already anticipated in ([11]) where it was obtained by considering a possible realization of space-time coordinates as differential operators in momentum space. The conclusion that one gets from (19-22) is that there is a modification of the quantum mechanical commutators which becomes relevant when the energy approaches its maximum value and in the limit one gets a classical phase space. This is a result that one could have anticipated from the consistency of the quantum mechanical uncertainty principle with the possibility to explore an arbitrarily small region in space-time while having a cutoff on the available energies.

IV. AN EXAMPLE WITH A MOMENTUM CUTOFF: DSR1

Another example of a non linear realization of Lorentz transformations corresponds to the choice of functions

\[
f = \frac{1}{2} \left[ (1 + \ell^2 p^2) e^{\ell E} - e^{-\ell E} \right]
\]

\[
g = e^{\ell E}
\]

The relation between the energy and momentum for a particle of mass \(m\) is given by

\[
(1 - \ell^2 p^2) e^{\ell E} + e^{-\ell E} = e^{\ell m} + e^{-\ell m}
\]

which is the dispersion relation of the model referred to as DSR1 [5]. One finds

\[
e^{\ell E} = \frac{\cosh(\ell m) + \sqrt{\cosh^2(\ell m) - (1 - \ell^2 p^2)}}{1 - \ell^2 p^2}
\]

for the energy as a function of momentum. One has in this case an upper bound on the momentum (\(p^2 < 1/\ell^2\)) instead of the energy. If one replaces the function \(g\) and its derivatives

\[
\varpartial_1 g = g = e^{\ell E} \quad \varpartial_2 g = 0
\]

into the general expression (7-10) for the modified quantum mechanical commutators one has

\[
[t, E] = i\hbar \frac{1}{D_1}
\]

\[
[t, p_i] = -i\hbar \ell p_i \frac{1}{D_1}
\]

\[
[x_i, E] = -i\hbar \ell \frac{2 e^{-\ell E} \ell E \varpartial_2 f}{D_1}
\]

\[
[x_i, p_j] = i\hbar e^{-\ell E} \left[ \delta_{ij} + \ell^2 p_i p_j \frac{2\ell E \varpartial_2 f}{D_1} \right]
\]

where

\[
D_1 = f + \ell E \left[ \varpartial_1 f - 2\ell^2 p^2 \varpartial_2 f \right]
\]

For the particular choice for \(f\) in (23)

\[
2\ell E \varpartial_2 f = e^{\ell E}
\]

and

\[
D_1 = \frac{1}{2} \left[ (1 - \ell^2 p^2) e^{\ell E} + e^{-\ell E} \right]
\]

If one uses the relation between energy and momentum (25), the right hand side in (34) reduces to \(\cosh(\ell m)\) and the phase space commutators become

\[
[t, E] = \frac{i\hbar}{\cosh(\ell m)}
\]

\[
[t, p_i] = -\ell p_i \frac{i\hbar}{\cosh(\ell m)}
\]

\[
[x_i, E] = -\ell p_i \frac{i\hbar}{\cosh(\ell m)}
\]

\[
[x_i, p_j] = i\hbar \left[ e^{-\ell E} \delta_{ij} + \ell^2 p_i p_j \frac{1}{\cosh(\ell m)} \right]
\]

We see from these expressions that when the mass \(m\) is much larger than the inverse of the length scale \(\ell\) all the commutators are (exponentially) small and a classical phase space is approached. This result suggests the possibility to relate the transition from the quantum behaviour at the microscopic level to the classical behavior at the macroscopic level with the modification of quantum mechanics induced by a modification of the relativity principle.
As a final remark one can consider the massless case where
\[ e^{\ell E} = \frac{1}{1 - \ell |\mathbf{p}|} \] (39)
and the modified commutators are
\[ [t, E] = i\hbar \] (40)
\[ [t, p_i] = -i\hbar p_i \] (41)
\[ [x_i, E] = -i\hbar p_i \] (42)
\[ [x_i, p_j] = i\hbar [(1 - \ell |\mathbf{p}|) \delta_{ij} + \ell^2 p_i p_j] \] (43)

In contrast to the case of a cutoff in the energy, when the momentum approaches its maximum value one has a non trivial limit for the commutators which differs from the canonical commutation relations.

V. CONCLUSIONS

The standard arguments leading to a minimum physical length beyond which it is not possible to go in the presence of gravity assume that in the flat space-time limit one has the standard QM uncertainty principle. If there is a remnant of gravity in the flat space-time limit as an invariant length and a modification of the QM commutators along the lines of the examples considered in this work then these arguments do not apply. In fact we have shown that in some cases one could have that instead of an obstruction to a localization in a physical system beyond a minimal length due to a modification of the quantum uncertainty by gravitational effects one could have the opposite situation where as a remnant of the QG theory the quantum uncertainty is diluted and the system approaches to the classical limit with no uncertainties in the high energy limit and/or for large masses.

If the modifications of QM suggested by the examples analyzed in this work apply to the flat space-time limit of the QG theory then our understanding of different physical systems should be reconsidered. The qualitative description of physical systems on a macroscopic scale, based on standard QM, could be altered in some cases. Also the discussion of black hole evaporation could be modified when one approaches the invariant length scale where quantum black holes become classical. Even the quantum mechanical aspects of the evolution of the Universe could differ from standard physics expectations. An study of some of the consequences of a modification of QM as well as a more systematic analysis of all possible QM commutators corresponding to the different nonlinear representations of Lorentz transformations in energy-momentum space and different ways to introduce the space-time sector are presently under investigation.

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